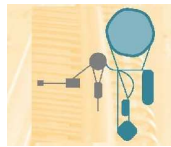


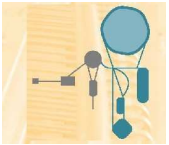
Studies of "electron heating" using MAD-X

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- Besides cooling, electron cooler acts as nonlinear optical element leading to excitation of resonances, and so to possibly resulting in emittance growth (effective "heating") → CELSIUS (V. Ziemann, TSL-note 98-43)
- Within FAIR, beams with high requirements to beam quality
→ electron cooling necessary especially for secondary beams in storage rings
- SchwerlonenSynchrotron (SIS) - 18:
 - has electron cooler
 - high currents leading to large space charge tune shift and spread
→ high probability that tune footprint crosses resonance
 - possibly experiment on this topic
- Extension of beambeam element by additional radial density profiles to describe interaction between particle beam and electron beam in MAD-X



- First rough estimate for influence of electron cooler from according linear tune shift
- Tune shift for electron beam with rectangular density may be written as

$$|\Delta\nu_z| = \frac{e}{8\pi^2\epsilon_0 m_0 c^2} \frac{\hat{\beta}_z e N'}{\beta^2 \gamma^3 b^2},$$

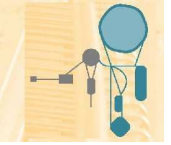
where $\hat{\beta}_z$ – beta function with $z = x, y$

N' – electron number in the cooler

β, γ – relativistic factors;

b – electron beam radius

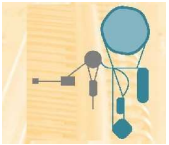
- **Stronger influence** for small energies → visible effect in SIS - 18 expected
- Includes electrostatic and magnetic interaction (2 parallel currents)



| | SIS - 18 ¹ | HESR ² |
|-----------------------|-----------------------|-------------------|
| pc/GeV | 0.15 (injection) | 1.5 (minimum) |
| γ | 1.01 | 1.89 |
| β | 0.15 | 0.85 |
| I_e/A | 0.3 | 0.2 |
| b/mm | 22 | 5 |
| l_{cooler}/m | 3 | 25 |
| $\hat{\beta}_x/m$ | 8 | 20 |
| $ \Delta\nu_x $ | 0.033 | 0.005 |

¹ B. Franczak, SIS Parameter List; L. Groening, GSI-dissertation 98-20.

² HESR Technical baseline report.



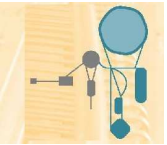
- Model taken from V. Ziemann, TSL-note 2004-60, is extension of model from E. Keil (CERN)
- Aim: Calculation of nonlinear tune shifts and footprints with electron beams with different density profiles and represented by a single kick
- using integral expressions for those:

$$\Delta\nu_x(J_x, J_y) = 2 \int_0^{2\pi} \frac{d\phi_x}{2\pi} \int_0^{2\pi} \frac{d\phi_y}{2\pi} \cos^2 \phi_x \frac{\tilde{S}(\sqrt{2J_x \cos^2 \phi_x + 2J_y \cos^2 \phi_y})}{\sqrt{2J_x \cos^2 \phi_x + 2J_y \cos^2 \phi_y}}$$

with action and angle variables, $J_{x/y} = \epsilon_{x/y}/2$ and $\phi_{x/y}$, of betatron motion

i.e. $x = \sqrt{2J_x \beta_x} \cos \phi_x$, y respectively,

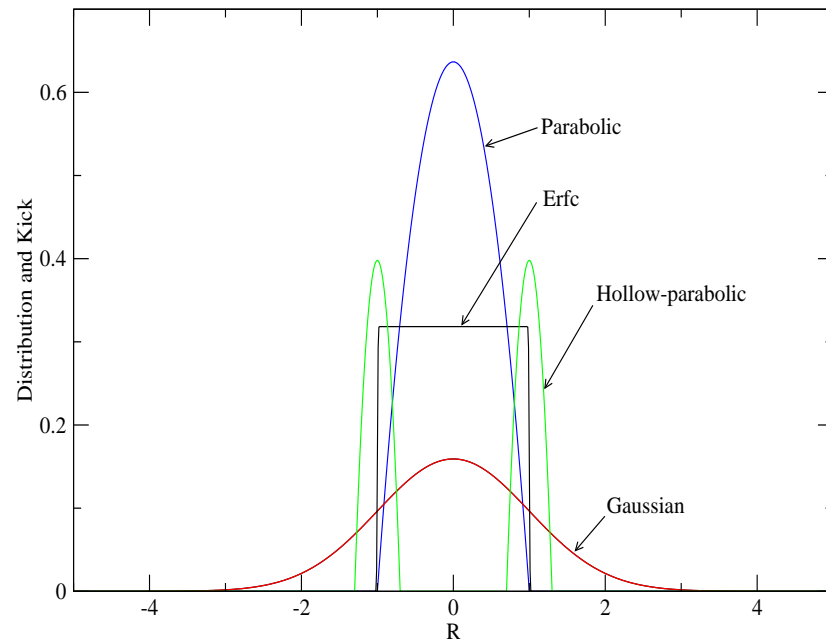
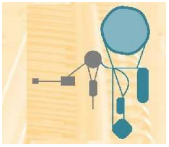
and the normalized electrostatic force function \tilde{S}



- Resonance strength from Fourier representation

$$\sqrt{2J} \cos \phi \tilde{S}(\sqrt{2J} \cos \phi) = J \sum_{n=0}^{\infty} A_n \cos n\phi$$

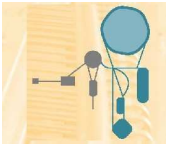
- Cooler kicks with several radial density distributions: **Gaussian**, **Parabolic**, **Error function (Flat top)**, and **Hollow parabolic**
- beambeam element with these shapes in **MAD-X** implemented excepted for parabolic profile



Charge normalized to 1, i.e.

$$2\pi \int_0^{\infty} dr r n_0(r) = 1$$

- Figure from **V. Ziemann**, TSL-note 2004-60, density profiles used: Gaussian, flat top (error function), hollow-parabolic, and parabolic

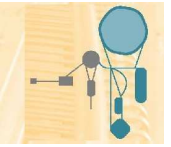


- Model system for tracking: rotational matrix + single beambeam kick

$$\begin{pmatrix} x_{n+1} \\ x'_{n+1} \end{pmatrix} = \begin{pmatrix} \cos 2\pi\nu_{0,x} & \hat{\beta}_{0,x} \sin 2\pi\nu_{0,x} \\ -\frac{1}{\hat{\beta}_{0,x}} \sin 2\pi\nu_{0,x} & \cos 2\pi\nu_{0,x} \end{pmatrix} \begin{pmatrix} x_n \\ x'_n + \Delta x'(x_n, y_n) \end{pmatrix}$$

with:

- phase space variables in i th step x_i, x'_i
- beta function in the cooler $\hat{\beta}_0$
- phase advance in the lattice $2\pi\nu_0$
- transverse momentum kick due to electron beam $\Delta x'(x_n, y_n)$
(thin lens approximation)

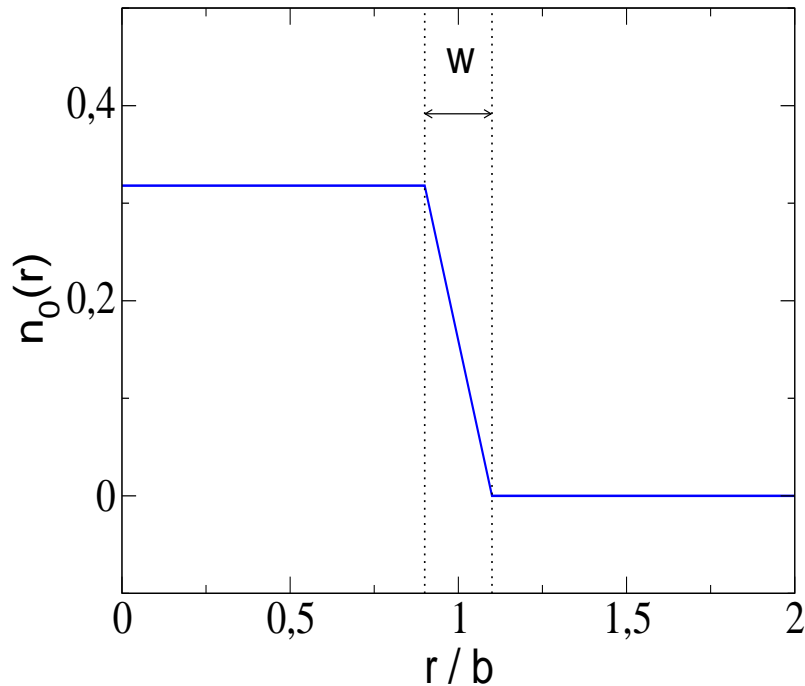
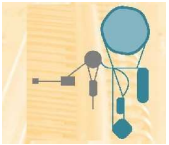


$$\Delta x'(x, y) = \frac{qq' N'}{2\pi\epsilon_0 m_0 c^2 \beta^2 \gamma^3} \frac{x}{R^2} \int_0^R dr r n_0(r)$$

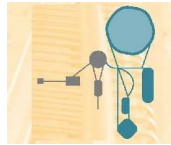
with

- q, q' – particle charge in considered beam (protons),
"colliding" beam (electrons)
- $N' = I_e L_{\text{cool}} / (\beta c)$ – number of electrons in the cooler
- $R = \sqrt{x^2 + y^2}$ – distance from the centre of the electron beam
- radial electron current density is normalized by (round electron beam)

$$\int_0^\infty dr r n_0(r) = 1$$



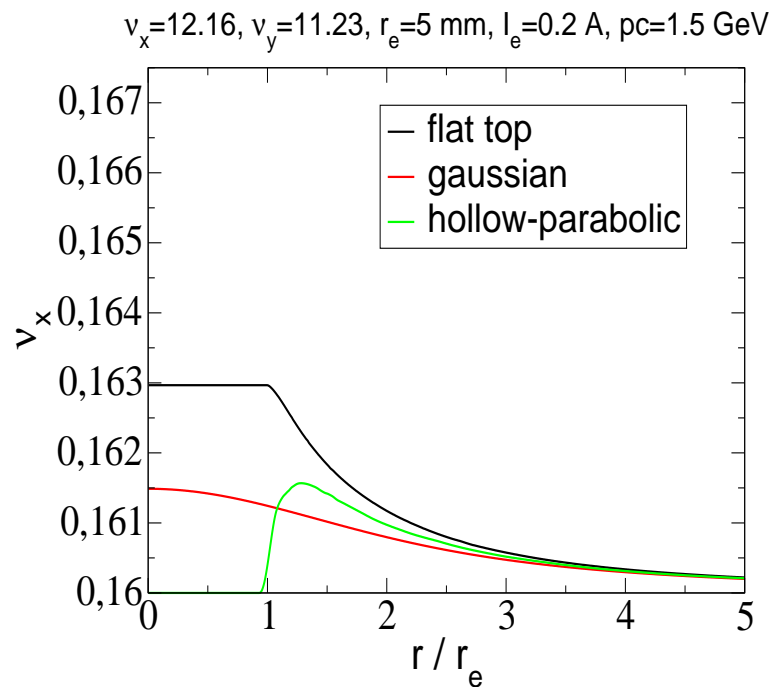
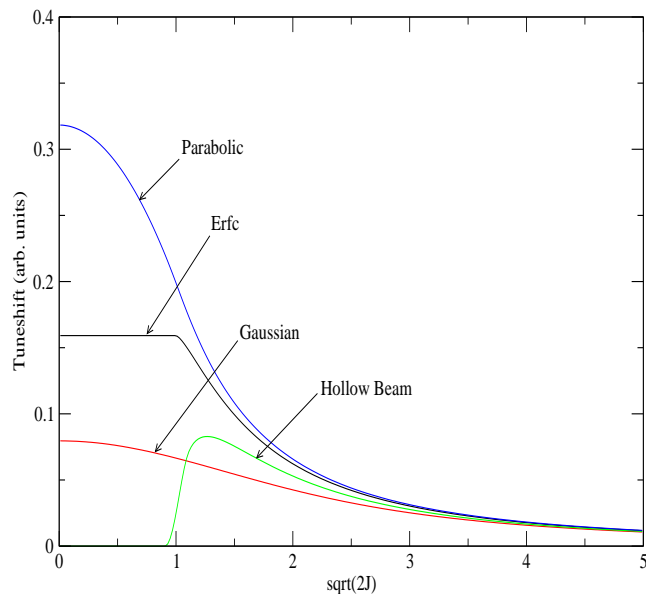
- beambeam kick elements in MAD-X with
 - Gaussian profile (initial)
 - Flat top profile (new)
 - Hollow-parabolic profile (new)
- Flat top profile for electron cooler important
- Here, trapezoidal profile used with beam radius b and width of edge layer w
- **Advantage:** analytical expressions for resulting force and maps



Non-linear tune shift due to the kick: **protons** \Rightarrow tune shift is positive

left: analytic, **right:** MAD-X, rotational matrix + single kick

HESR conditions: $(\nu_x = 12.16, \nu_y = 11.23)^1$, $I_e = 0.2$ A, $pc = 1.5$ GeV

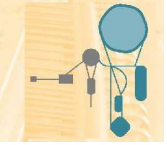


$$\frac{\hat{\beta}_z e^2 N'}{4\pi\epsilon_0 m_0 c^2 \beta^2 \gamma^3} \frac{1}{2\pi b^2}$$

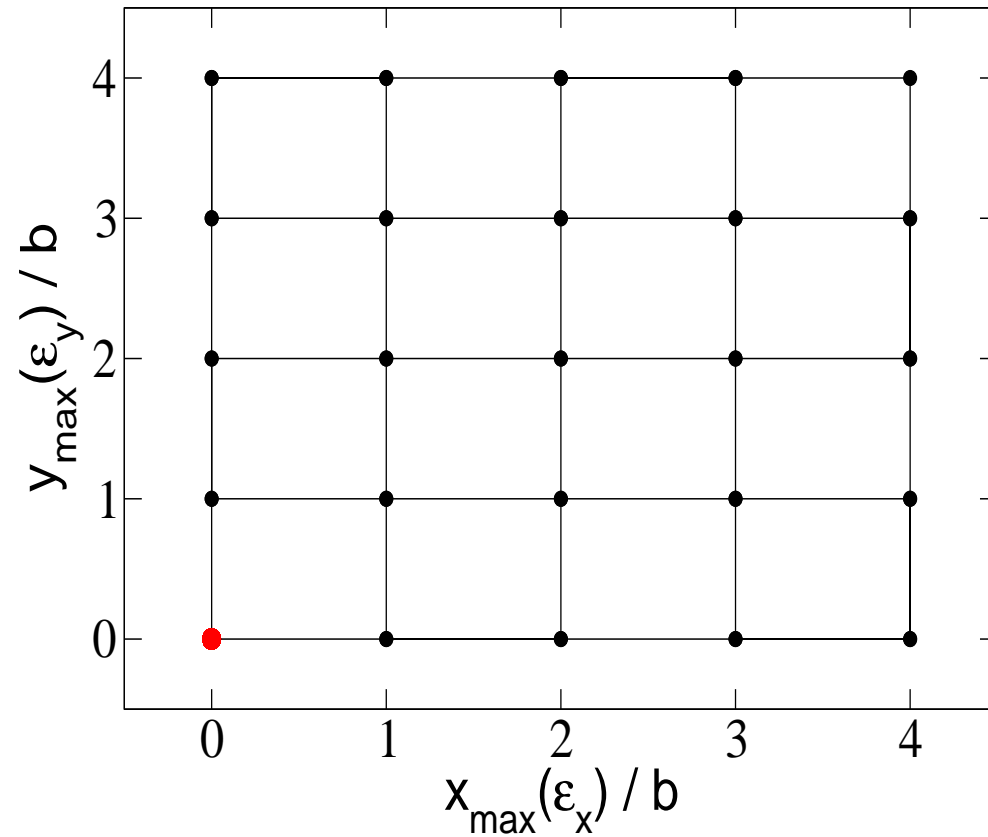
$$= \frac{1}{2\pi} \text{ (left)}$$

$$= 2.98 \cdot 10^{-3} \text{ (right)}$$

¹ Yu. Senichev, "The advanced HESR lattice for the stochastic cooling", FZ Jülich 2006

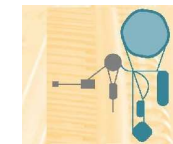


Tune footprint

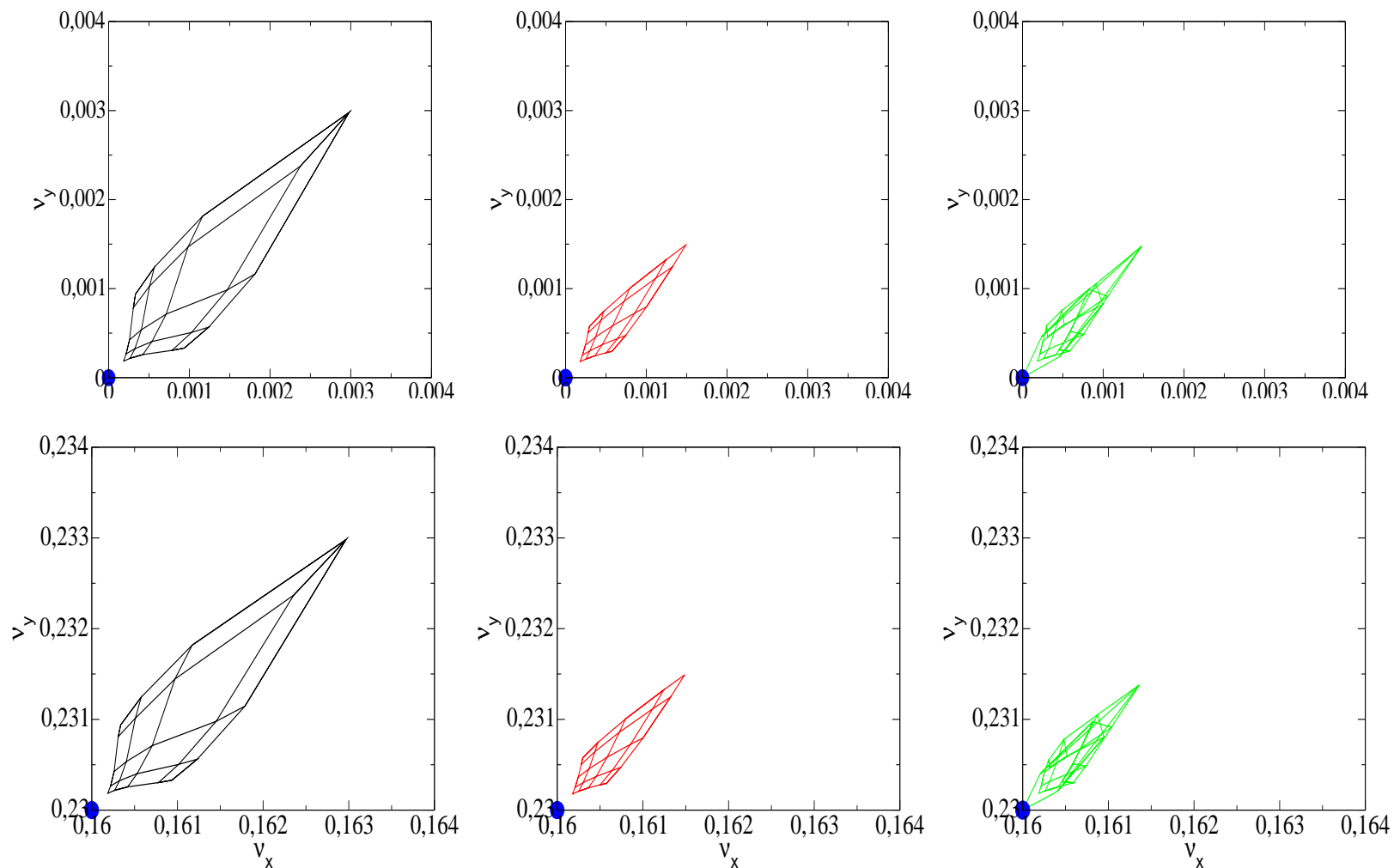


- Initial distribution chosen to make shape and structure of tune spreads better visible

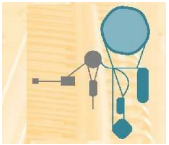
- $z_{\max} = \sqrt{\epsilon_z \hat{\beta}_z}$, $z = x, y$



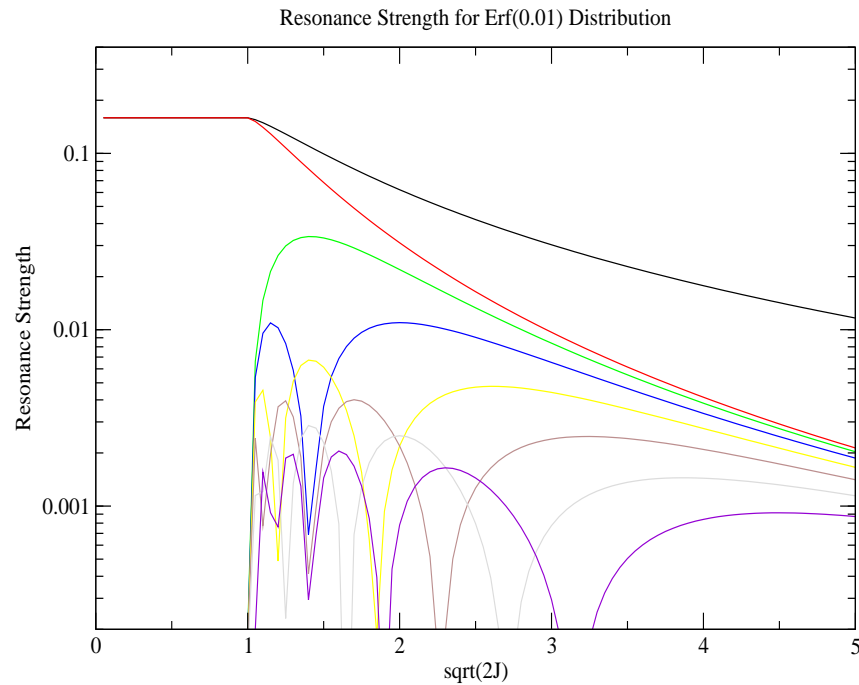
Tune footprints **up:** analytic, **down:** MAD-X, rotational matrix + single kick



Protons with $\nu_x = 12.16$, $\nu_y = 11.23$, $r_e = 5$ mm, $I_e = 0.2$ A, $pc = 1.5$ GeV

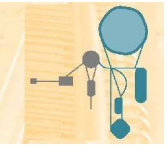


Use resonance strengths from analytic model to find relevant resonances.



- Flat top electron current profile (error function)
- Only resonances with even order appear
- Resonance strength strongly decreases with increasing order

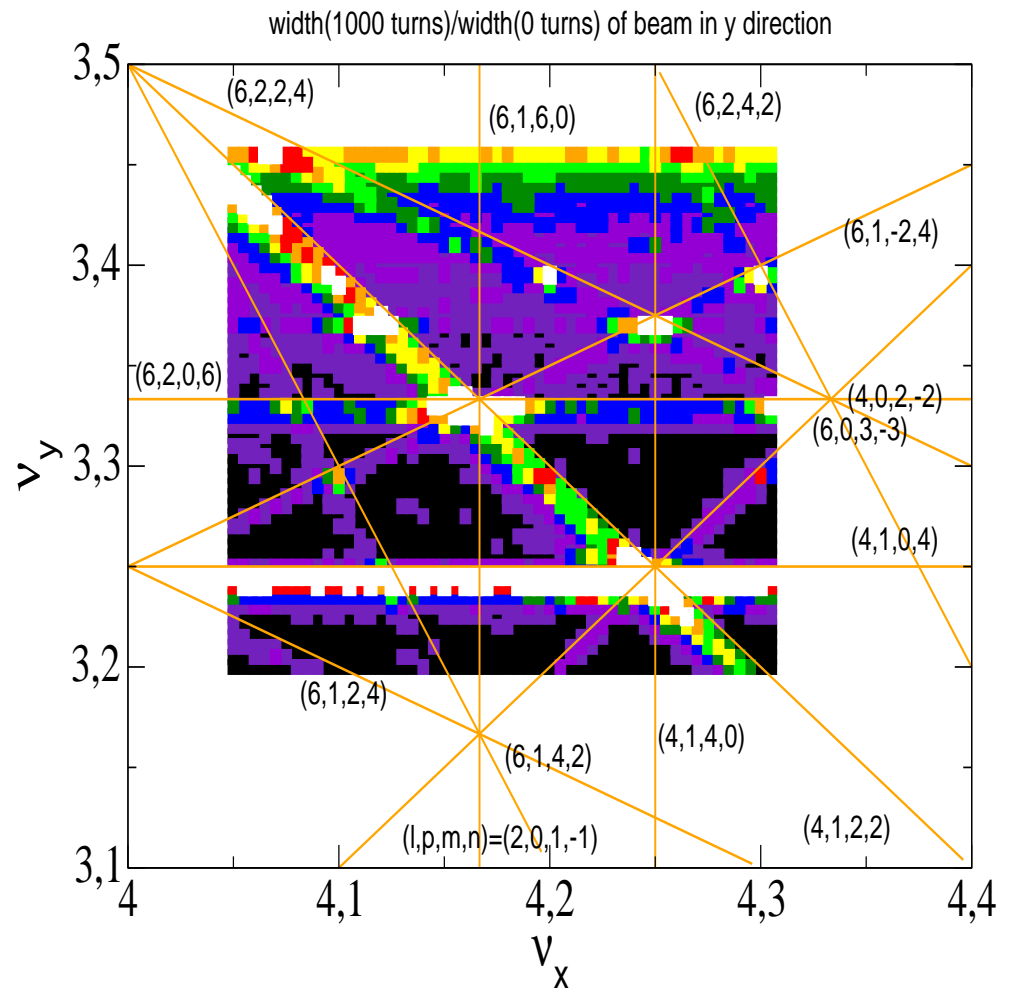
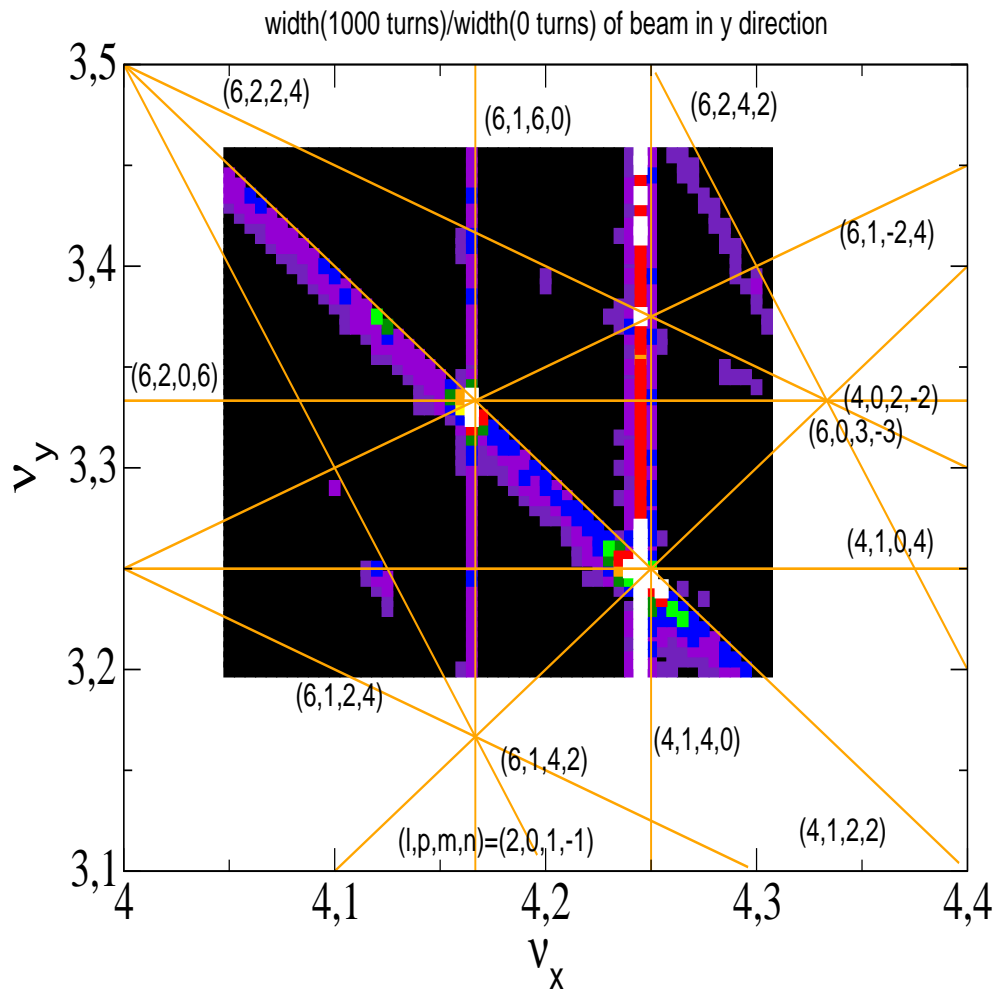
Figure from V. Ziemann, TSL-note 2004-60

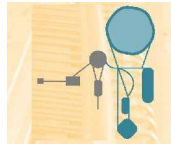


Relative beam width $w_{rel} = w_{fin}/w_{ini}$ depending on the tune

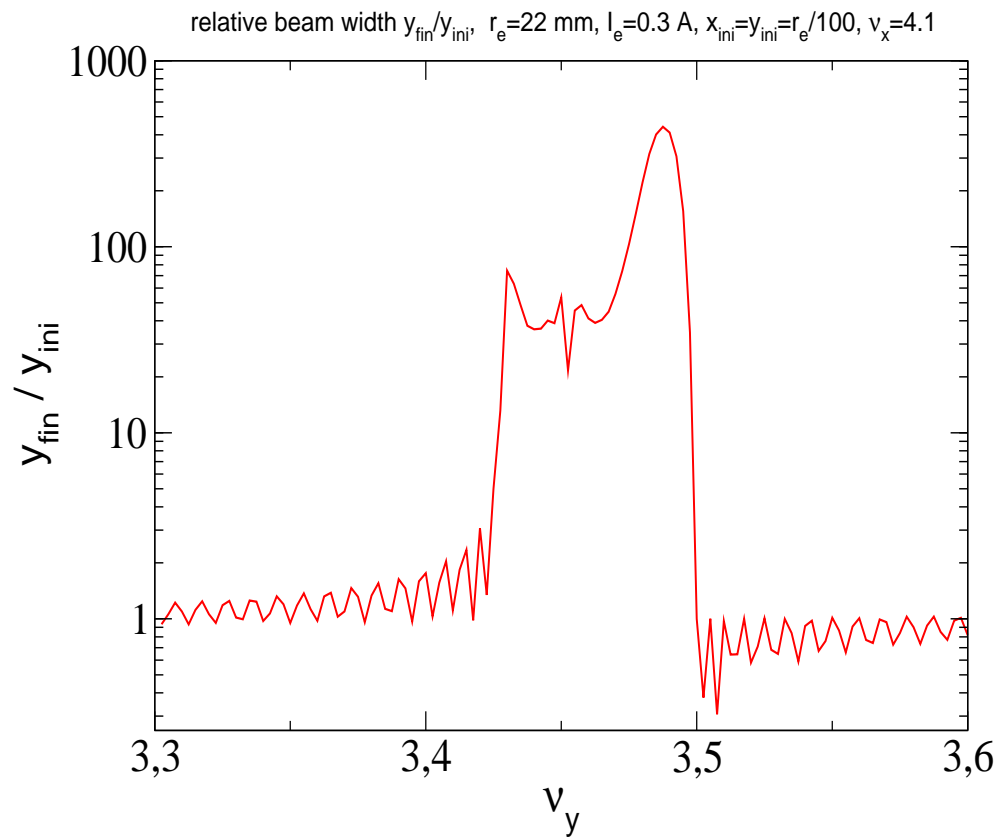
x direction

y direction

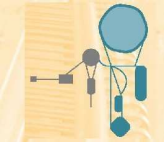




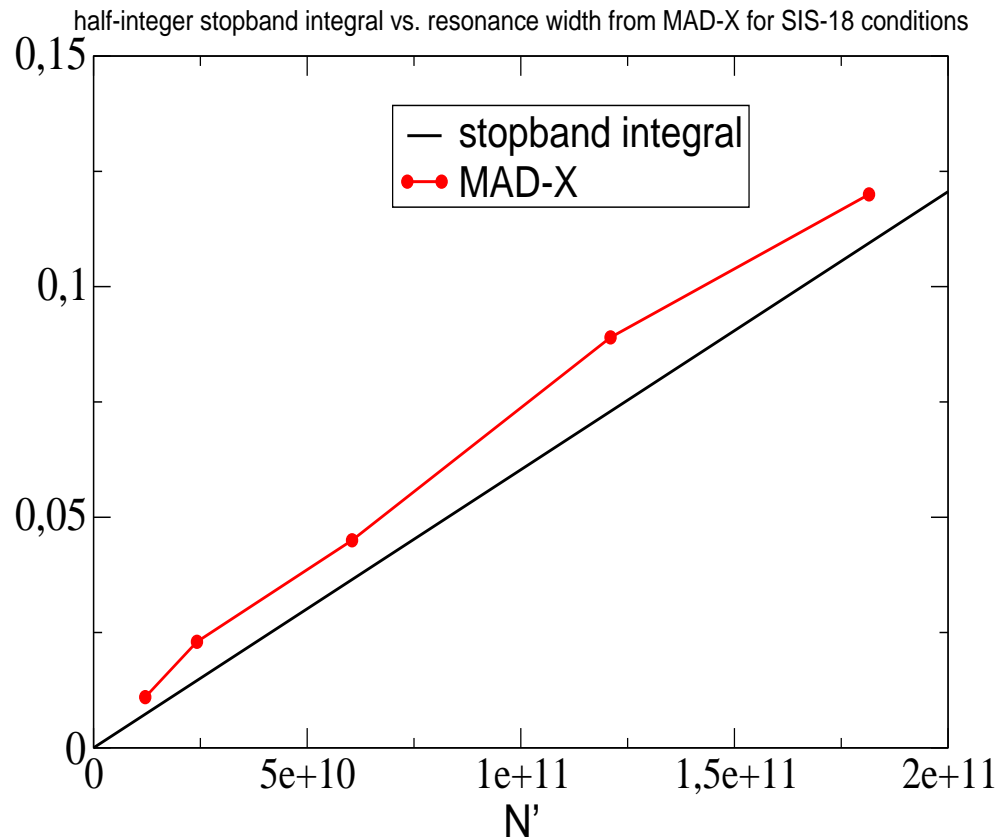
Relative beam width in y direction depending on the tune, $\nu_x = 4.1$



- $w_{\text{ini}} = b/100 \ll b$
 → only quadrupole error due to electron beam.
- Near half integer tune, influence of the quadrupole error due to electron beam becomes significant.
- Increase of the beam width occurs only **below** half integer tune.



Width of half integer resonance vs. half-integer stopband integral



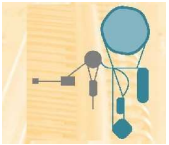
- half-integer stopband integral

$$J_p = \frac{1}{2\pi} \oint \hat{\beta} k(s) e^{-ip\phi} ds$$

- denotes the tune range, where

$$\left| \frac{\Delta \hat{\beta}}{\hat{\beta}} \right| = 1,$$

i.e. beam size increases by a factor $\sqrt{2}$



- beambeam element has been modified to make it usable for the description of an electron cooler
- Modified element has been tested, results have been compared to those of analytic model with good agreement
- First application: Calculation of resonances driven by the electron cooler in SIS - 18
 - major resonances could be identified
 - results seem reasonable
- Possible experiments being under discussion:
 - measurement of resonances
 - Measurement of non-linear tunes, i.e. depending on betatron amplitude